

Physics 850: Soft Condensed Matter Physics, Fall04
A.D. Dinsmore
Lecture 5: Polymers

Polymers

(Kleiman et al)
(Rubenstein & Colby Polymer Physics)

Polymer = "molecule of high molecular mass composed of many small structural units connected by strong covalent bonds"

Types:

homopolymers - units are identical

copolymers - > 1 type of monomer

random copolymers - units vary randomly

block copolymers - units are identical in sections

(di - 2 blocks

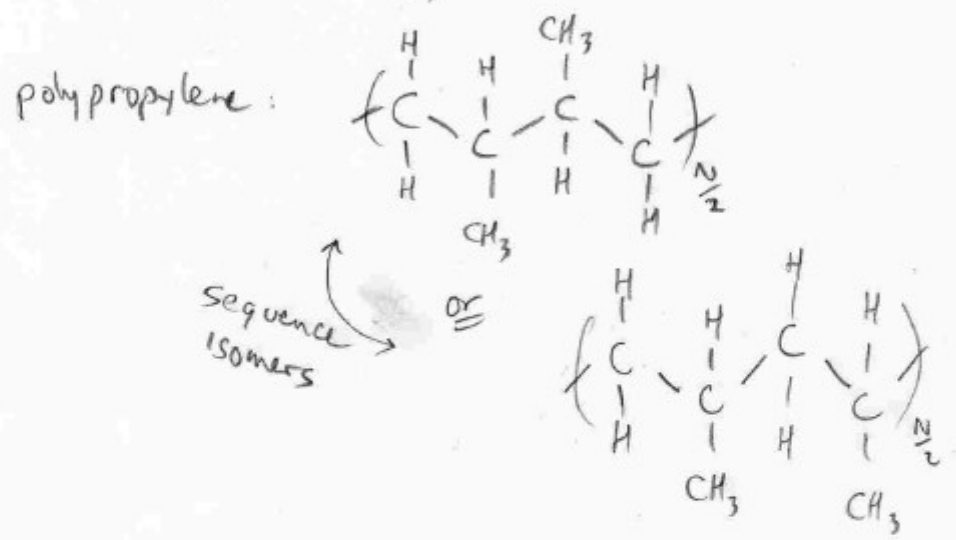
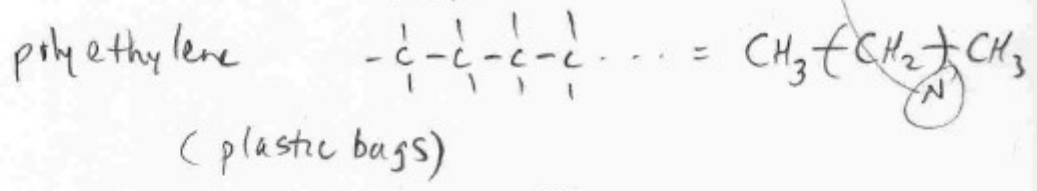
tri - 3 blocks)

linear: each unit has 2 bonds (except ends)

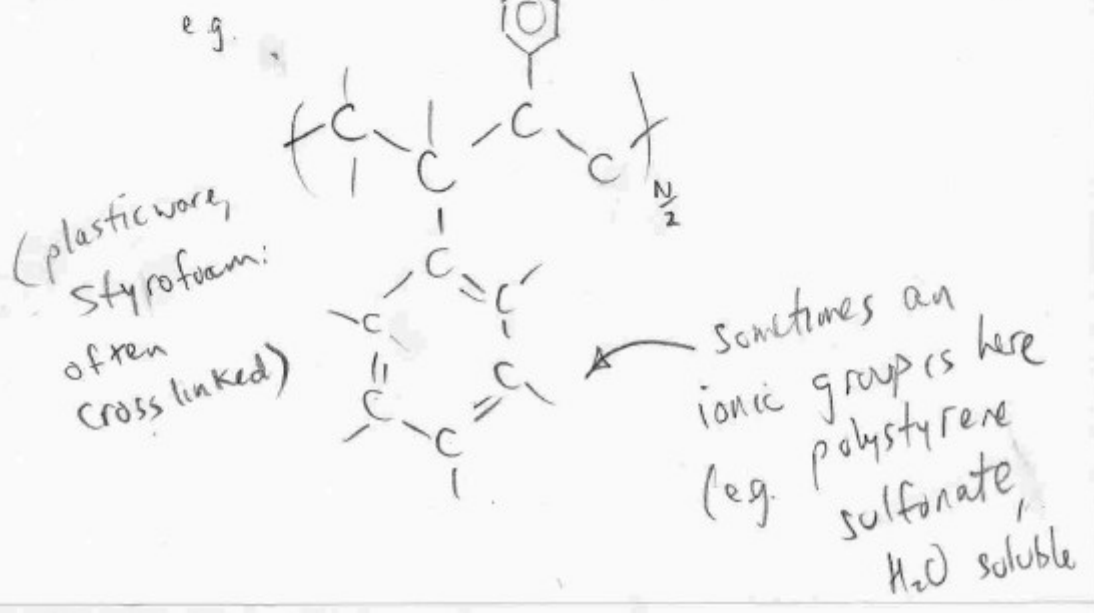
branched, ring, star, dendrimer, comb, ladder, ...

Degree of polymerization

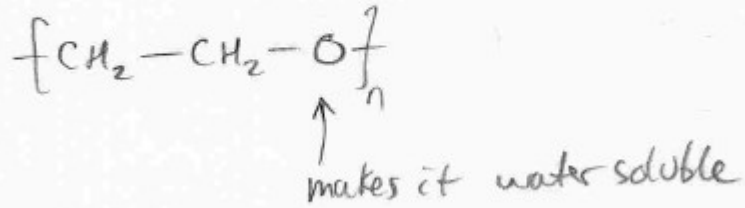
Examples of linear homopolymers ≈ 1.5 A



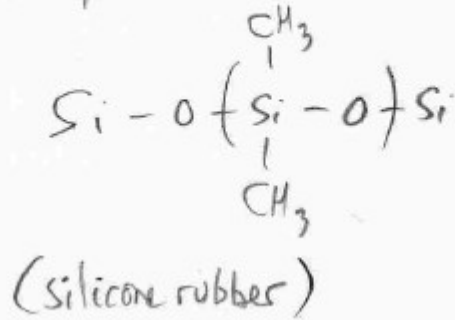
polystyrene: (replace CH₃ of polypropylene with C ring)



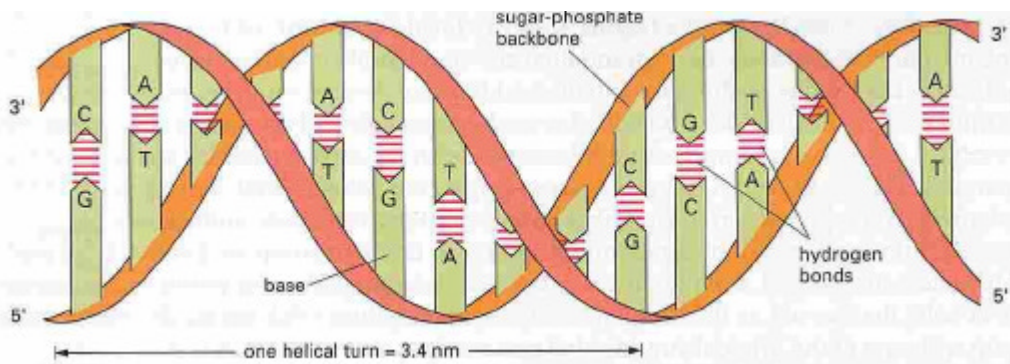
(PEG)
 polyethylene oxide, polyethylene glycol, poly-oxyethylene.



polydimethyl siloxane (PDMS)

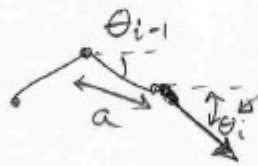


Double-stranded DNA: a (pseudo) random copolymer with 4 bases: A, G, C, and T. Kuhn length ~ 100 nm (much larger than polyethylene's 1.4 nm).



A simple model: Ideal (Gaussian) chain (Freely jointed chain)

- Assume polymer is composed of N segments of length a , which do not interact and are independent (except by being connected)



θ_i is independent of θ_{i-1}

• Each junction is a freely pivoting hinge

• \hat{t}_i is direction of i^{th} segment

→ contour length = Na

→ Average lengths:



Let this end sit at origin

this end at $(x, y, z) = \vec{r}$

$$\vec{r} = \sum_{i=1}^N a \hat{t}_i$$

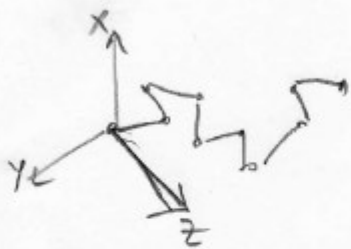
$$\langle \vec{r} \rangle = aN \langle \hat{t} \rangle = 0 \text{ because } \theta_i \text{ is random}$$

$$\langle r^2 \rangle = \langle (a \sum_i \hat{t}_i) \cdot (a \sum_j \hat{t}_j) \rangle = a^2 \langle \sum_{i,j=1}^N \hat{t}_i \cdot \hat{t}_j \rangle$$

\leftarrow only N terms with $i=j$ contribute.

$$\text{so } \boxed{\langle r^2 \rangle = Na^2}$$

Central-limit - Theorem (cf. Kleman §15.1)
 or Random-walk view:



x = end-to-end length of polymer along x -axis
 y = " " " " y -axis
 z = " " " " z -axis

$$P(x, y, z) = (\text{const.}) e^{-\frac{(x-\langle x \rangle)^2}{2\sigma_x^2}} e^{-\frac{(y-\langle y \rangle)^2}{2\sigma_y^2}} e^{-\frac{(z-\langle z \rangle)^2}{2\sigma_z^2}}$$

$\rightarrow \langle z \rangle = \langle x \rangle = \langle y \rangle = 0$ (polymer equally likely to fluctuate to L or R)
 so $\langle (x-\langle x \rangle)^2 \rangle = \langle x^2 \rangle$
 and $\sigma_x^2 = \langle x^2 \rangle$

$$\rightarrow \langle x^2 \rangle + \langle y^2 \rangle + \langle z^2 \rangle = \langle r^2 \rangle \rightarrow \langle x^2 \rangle = \frac{\langle r^2 \rangle}{3} = \langle y^2 \rangle$$

$$P(r) = (\text{const.}) e^{-\frac{r^2}{2\langle r^2 \rangle}} \quad \text{where } r^2 = x^2 + y^2 + z^2$$

$$= (\text{const.}) e^{-\frac{3r^2}{2\langle r^2 \rangle}}$$

Since $P(r) \propto e^{-\frac{F(r)}{kT}} \rightarrow F(r) = F_0 + \frac{3k_B T r^2}{2\langle r^2 \rangle}$

$$\rightarrow \langle r^2 \rangle = a^2 N$$

$$\text{so } F(r) = F_0 + \frac{3k_B T r^2}{2a^2 N}$$

force exerted by polymer = $-\frac{\partial F(r)}{\partial r} = -\frac{3k_B T}{a^2 N} r$ — spring constant

Corrections to Ideal chain model

- 1) What is the random walk length?
(Kuhn length, b)

linear hydrocarbon polymers rotate by
alternation between *trans* & *gauche*
(Kleman Fig 15.1)

Polyethylene:

C-C bond length $\approx 1.5 \text{ \AA}$

Kuhn length, $b \approx 14 \text{ \AA}$ because of correlations

Trans-gauche variations in structure give rise to bends – see Kleman

DNA (double helix)

$b \approx 100 \text{ nm}$ (70x larger than PE)

- 2) Excluded volume: 2 chain segments cannot overlap

Result: $\sqrt{\langle r^2 \rangle} = a \cdot N^{3/5}$ (not $\frac{1}{2}$, as
in ideal
model)

(Flory theory)